

Progress of high performance simulations of an accretion disk surrounding a supermassive black hole

Fabian Klein ¹, Rainer Spurzem ¹²³, Andreas Just ¹, Rolf Kuiper⁴

¹Astronomisches Rechen-Institut, Zentrum für Astronomie Heidelberg, Universität Heidelberg, Mönchhof-Straße 12-14, D-69120 Heidelberg, Germany

²National Astronomical Observatories of China, Chinese Academy of Sciences, NAOC/CAS, 20A Datun Rd., Chaoyang District, Beijing 100012, China

³The Kavli Institute for Astronomy and Astrophysics at Peking University

⁴Universität Tübingen, Institut für Astronomie und Astrophysik, Abt. Computational Physics, Auf der Morgenstelle 10, 72076 Tübingen, German

31st of May 2017

Table of Contents

- Introduction
- Modeling and Initial conditions
- First results
- Benchmarks
- Conclusions

STARDISK

Simulations of an SMBH in an AGN with the surrounding star-cluster (Nbody, see [Just et al., 2012, Kennedy et al., 2016])
Accretion disk of black hole uses modified Shakura-Sunyaev model
(No Simulation)
Result: Accretion rate increased by presence of Accretion disk

My contribution

Hydrodynamical simulation of the disk instead of just enforcing model.
Interactions of Disk with Star-cluster, mutual feedback
Calculation of star-crossings (heating expected)



Figure: Artist's impression of a super massive black hole with an accretion disk, Image credit: NASA/JPL-Caltech

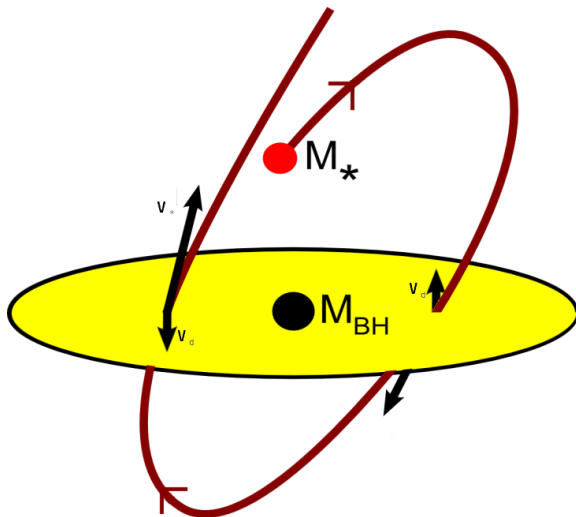


Figure: Figure illustrating the STARDISK situation, Drawing by Gareth F. Kennedy, modifications by Bekdaulet Shukirgaliyev

Paper I

[Just et al., 2012]

- Stationary Keplerian rotating disk
- Disk has constant thickness
- Star accretion rate enhanced by STARDISK interaction compared to pure stellar dynamics

Paper II

[Kennedy et al., 2016]

- Linear thickness in R compared to constant thickness
- Detailed orbit analysis

Foundations

[Novikov and Thorne, 1973, Shakura and Sunyaev, 1973]

- Assume turbulent disk with turbulent velocity v_T
- model effect of magnetic field H by viscosity tensor σ
- Amount of turbulence (compare to speed of sound c_S^2) determines angular momentum transport by viscosity

$$\alpha = \frac{v_T}{c_S} + \frac{H^2}{4\pi\rho c_S^2} \quad (1)$$

Details

- Generally $\alpha < 1$
- Different regions by p_{rad} in comparison to p_{gas}
- Process of opacity Thompson scattering, free-free
- Only $\sigma_{R\varphi} = \alpha\rho c_S$ relevant

Features

- Predicts radial and to some degree z profiles for density, temperature etc.
- Is able to partly explain AGN spectra
- Is theoretically implied to work for non-self-gravitating disks with local MHD turbulences [Balbus and Papaloizou, 1999]
- Radiation transport, radiation pressure, viscosity and hydrodynamics necessary
- Accretion flow and rate determined by α parameter

Governing equations

- Need: modified Navier-stokes equations of hydrodynamics
- Ideal equation of state is used to close the system
 $p = \rho \sigma_{\text{rms}}^2 = \rho \frac{k_B}{\mu_{\text{mol}} m_H} T$, also used to initialise temperature
- Radiation transport included in Flux-Limited-diffusion limit
- Currently: Region such that self-gravity can be neglected

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0 \quad (1)$$

$$\frac{\partial \mathbf{m}}{\partial t} + \nabla (\mathbf{m} \cdot \mathbf{v}) + \nabla p - \nabla \sigma = \rho \mathbf{g} + \rho \frac{\kappa}{c} \mathbf{F} \quad (2)$$

$$\frac{\partial E}{\partial t} + \nabla (E \mathbf{v}) + \nabla (p \mathbf{v}) - \nabla \cdot (\sigma \mathbf{v}) = \mathbf{m} \cdot \mathbf{g} + \rho \frac{\kappa}{c} \mathbf{v} \cdot \mathbf{F} + \rho c \kappa (a T^4 - E_R) \quad (3)$$

$$\frac{\partial E_R}{\partial t} + \nabla \cdot (E_R \mathbf{v}) = -\nabla \cdot \mathbf{F} - \rho c \kappa (a T^4 - E_R) \quad (4)$$

Equilibrium conditions derived from force Equilibrium [Lodato, 2008]. $v_r = v_\vartheta = 0.0$.

- Employ finite volume grid based code PLUTO
[Mignone et al., 2007]
- Use self-gravity and radiation transport modules as well more PLUTO additions developed by Rolf Kuiper
[Kuiper et al., 2010]
- Make use of heavy MPI parallelisation provided by PLUTO and the modules

Units

Chosse M31 as given in [Kennedy et al., 2016]

$$M_{\text{BH}} = 1.5 \times 10^8 M_{\odot}$$

$$M_{\text{Disk}} = 1.5 \times 10^7 M_{\odot}$$

Grid

- Spherical coordinates r, θ, φ
- Logarithmic grid in r , uniform in θ and φ

$$r \in [0.11 \text{ pc}, 2.95 \text{ pc}], \theta \in [0.472 \pi, 0.527 \pi], \varphi \in [0, 2\pi]$$

Grid and Units

Grid

- Spherical coordinates r, θ, φ
- Logarithmic grid in r , uniform in θ and φ

$$r \in [0.11 \text{ pc}, 2.95 \text{ pc}], \theta \in [0.472 \pi, 0.527 \pi], \varphi \in [0, 2\pi]$$

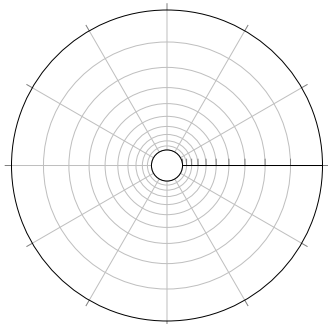


Figure: View from above on r, φ grid

Parameter collection

- $R_{\text{SW}} = \frac{2GM_{\text{BH}}}{c^2} \approx 3.589 \times 10^{-5}$ pc Schwarzschild radius
- $M_{\text{Cl}} = 1.5 \times 10^9 M_{\odot}$ Nuclear Star Cluster Mass
- $M_{\text{BH}} = 1.5 \times 10^8 M_{\odot}$ Black Hole Mass
- $M_{\text{d}} = 1.5 \times 10^7 M_{\odot}$ Total Disk Mass
- $R_{\text{d}} = 25$ pc Outermost radius of the disk
- $R_{\text{sg}} = \left(\frac{1}{2-p} \frac{M_{\text{BH}}}{M_{\text{d}}} h_z R_{\text{d}}^{2-p} \right)^{\frac{1}{3-p}} \approx 2.95$ pc Self-gravity radius, starting from which self-gravity is important
- $h_z = 1.0 \times 10^{-3} R_{\text{d}} = 0.025$ pc Scale height
- $h(R) = \begin{cases} \frac{h_z}{R_{\text{sg}}} R & \text{if } R < R_{\text{sg}} \\ h_z & \text{else} \end{cases}$ Linearly varying scale height in BH gravity dominated area

$R = r \sin \theta$ cylindrical radius, $T_{\text{K, inner}} = 2.93 \times 10^2$ yrs,
 $T_{\text{K, outer}} = 3.88 \times 10^4$ yrs

Current initial conditions

First Goal: Equilibrium initial conditions \Rightarrow Stationary disk
Derive stationary state from system of equations in 3 spatial dimensions [Lodato, 2008, Just et al., 2012, Kennedy et al., 2016].

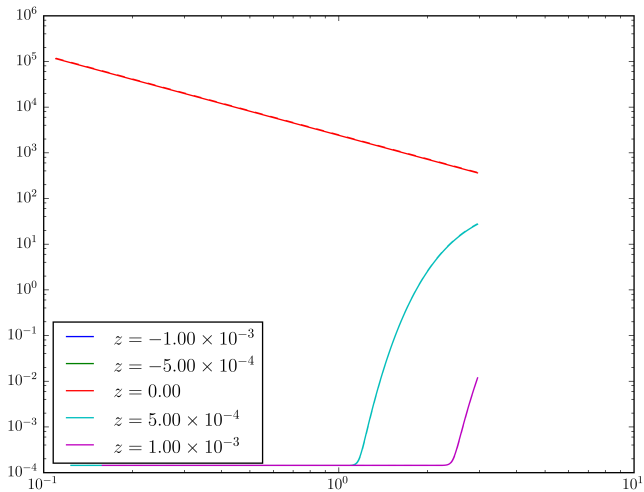
$$\rho(r, \theta) = \rho_{R_{0\rho}} \left(\frac{r}{R_{0\rho}} \right)^{\beta_\rho} \exp \left(-\frac{\cos^2 \theta}{2a_{R_{0as}}^2} \left(\frac{r}{R_{0as}} \right)^{-2y_{as}} \right)$$

$$v_\varphi = r \sin \theta \sqrt{\frac{GM_{\text{BH}}}{r^3}} \sqrt{\left(1 + \frac{\beta_\rho + \beta_T}{\gamma_{\text{gas}}} \frac{h^2}{r^2 \sin^2 \theta} \right)}$$

$$c_s = \frac{h}{r \sin \theta} v_\varphi \quad p = \frac{1}{\gamma_{\text{gas}}} c_s^2 \rho$$

We always choose $y_{as} = 0$, $a = \frac{h}{R}$. Also, I infer $\rho_{\text{min}} = 1.0 \times 10^{-23} \text{ g cm}^{-3}$ (6H cm^{-3}), Currently freezing in region with lower density initially.

Current initial conditions



Boundary conditions

r

- Beginning: Outflow no inflow, End Outflow no inflow
- Zero gradient in v_θ
- v_φ : Modified Keplerian gradient towards computational domain

θ

- Beginning: Outflow no inflow, End Outflow no inflow
- Zero gradient in v_r
- v_φ : Modified Keplerian gradient towards computational domain

φ

Periodic for end and beginning

Modified Keplerian gradient towards computational domain(example r):

Loop over all cells of the boundary zone, 2 cells in r extended over entire θ and φ .

$$v'_{\varphi} = v_{\varphi} \sqrt{\frac{r_{\text{outermost}}}{r_{\text{current}}}} \quad (5)$$

Limit speed to modified Keplerian speed

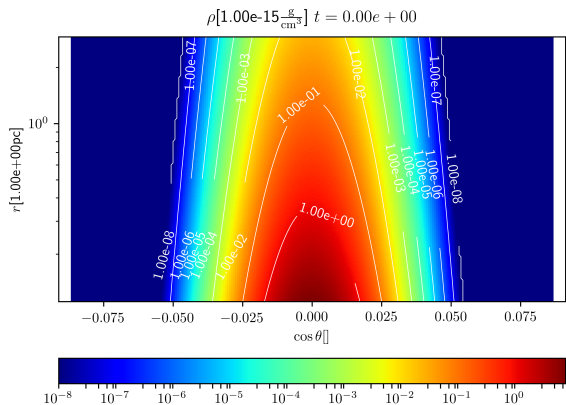
$$v''_{\varphi} = \text{Min}(v'_{\varphi}, v_{\text{K, modified}}) \quad (6)$$

2D test without radiation pressure and self-gravity deactivated

Start with 2D test, $\alpha = 5.0 \times 10^{-2}$, everything apart from self-gravity activated.

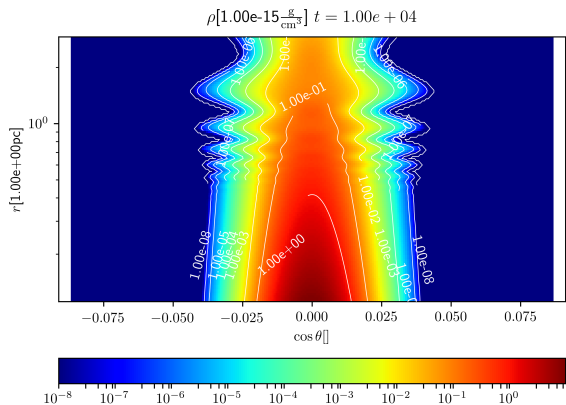
Runs fine up to at least 5×10^5 yrs

2D test without radiation pressure and self-gravity deactivated



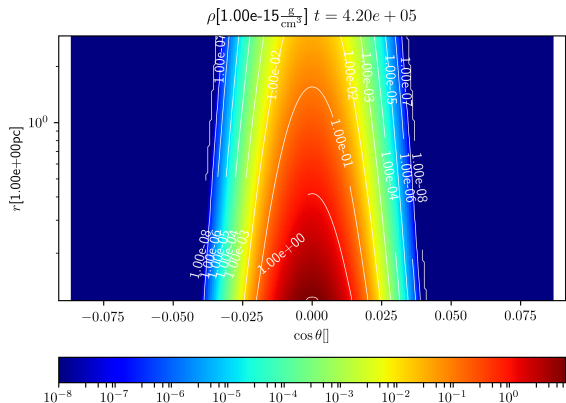
Initial conditons

2D test without radiation pressure and self-gravity deactivated



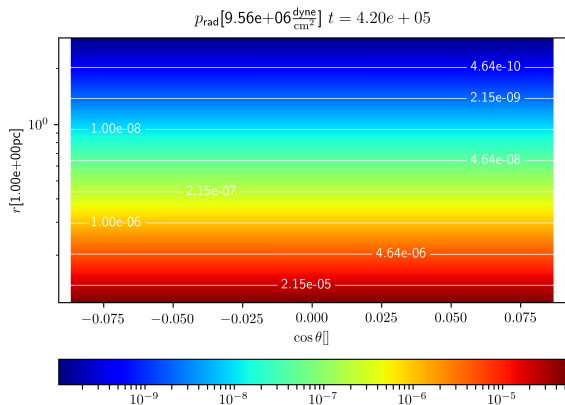
Initial oscillations

2D test without radiation pressure and self-gravity deactivated



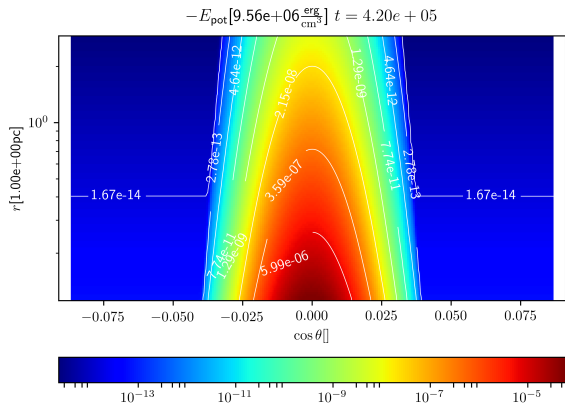
Equilibrium

2D test without radiation pressure and self-gravity deactivated



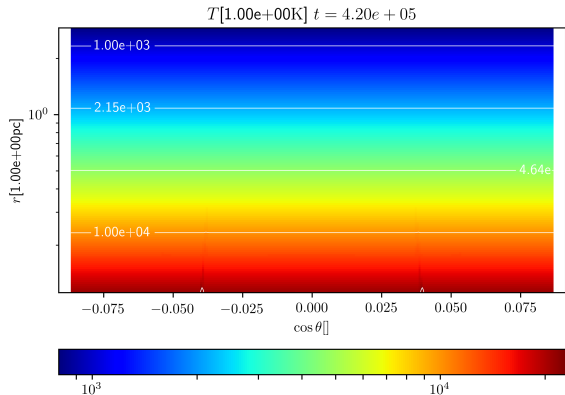
Radiation pressure dominates

2D test without radiation pressure and self-gravity deactivated



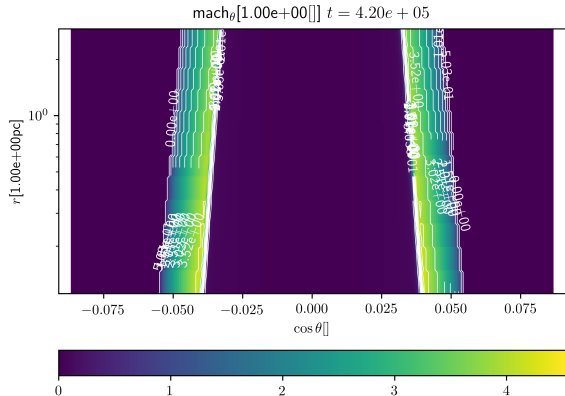
Greater than potential energy

2D test without radiation pressure and self-gravity deactivated



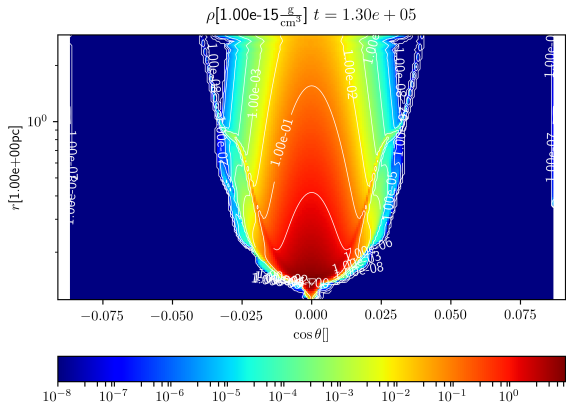
Temperature unchanged

2D test without radiation pressure and self-gravity deactivated



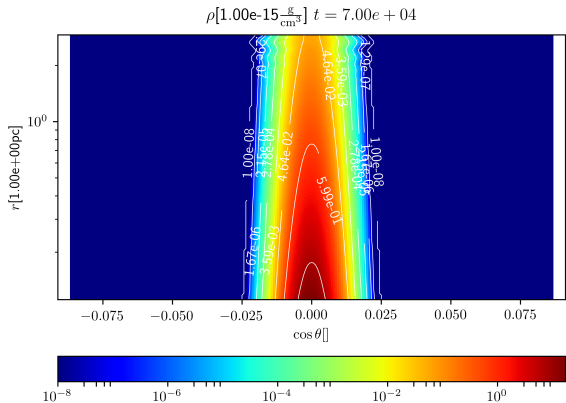
Shock in low density area

Problem: Too high radiation pressure causes explosion



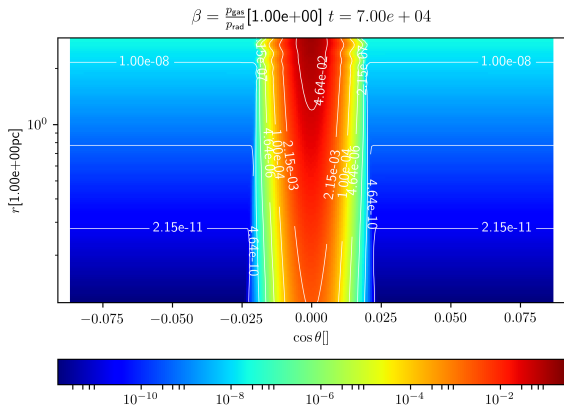
Explosion from much higher pressure in system

Idea: Reduce initial pressure by factor of 3



Results in thinner disk

Idea: Reduce initial pressure by factor of 3



Somewhat better, but still direct crash

Determination of critical temperature

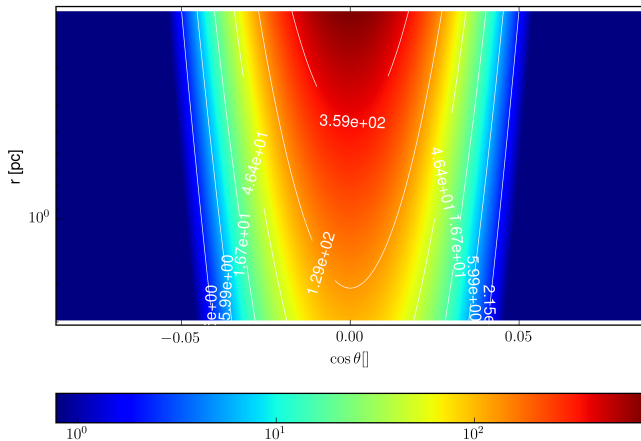
Assume density profile and ideal eos, calculate Temperature

$$p = \rho \frac{k_B T}{\mu m_H} \quad (7)$$

$$\beta = \frac{\rho_{\text{gas}}}{\rho_{\text{rad}}} = \rho \frac{k_B T}{\mu m_H} \frac{3}{a_{\text{rad}} T^4} \quad (8)$$

$$T_{\text{crit}} = \sqrt[3]{\frac{1}{\beta} \rho \frac{3k_B}{\mu m_H a_{\text{rad}}}} \quad (9)$$

Determination of critical temperature



$$\frac{\partial p_{\text{rad}}}{\partial r} = -\rho \frac{GM}{r^2} \quad (7)$$

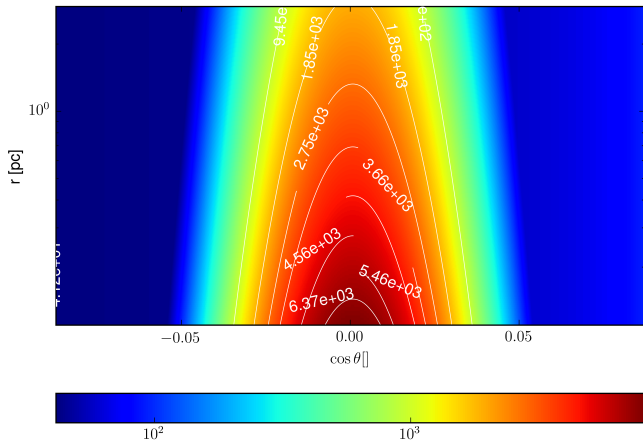
$$4a_{\text{rad}} T^3 \frac{\partial T}{\partial r} = -\rho \frac{GM}{r^2} \quad (8)$$

$$T(r) = \sqrt[4]{\frac{a_{\text{rad}} c_1 r + \rho GM}{a_{\text{rad}} r}} \quad (9)$$

$$c_1 = \frac{a_{\text{rad}} T_0^4 r_0 - GM\rho(r_0, \theta)}{a_{\text{rad}} r_0} \quad (10)$$

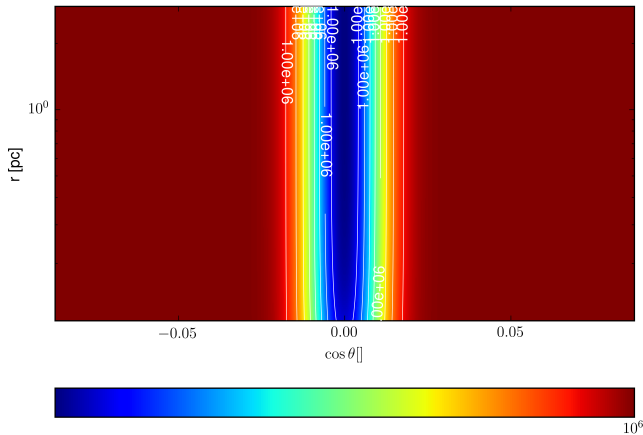
Assume $c_1 = 0.0$, Trying to get it via boundary, some negative values

Radiation pressure equilibrium



with $c_1 = 0$

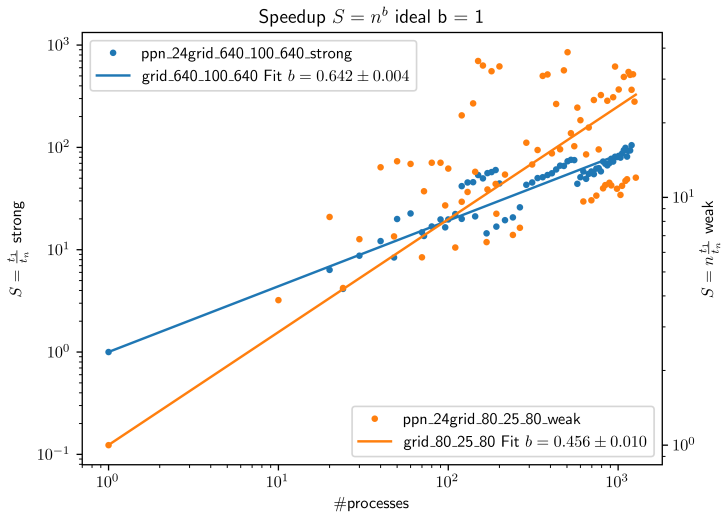
Radiation pressure equilibrium



with c_1 chosen such that $T(0.11 \text{ pc}) = 1.0 \times 10^6 \text{ K}$

- Disk goes to equilibrium, stable for long periods of time
- Radiation pressure very high, theoretically cannot be ignored.
- Causes explosion, no equilibrium any more
- Self-Gravity needs to be tested

Benchmarks I



Strong scaling OK, problem with weak scaling

Further goals

- Slowly progress to 3D simulation
- Do some benchmarks
- Extend to region where self-gravity is important
- Introduce heating due to star-crossings and interactions
- Determine spectra from radiation transport results
- Far goal: Include Hydrodynamical simulation in N-Body simulation

Thank you for your attention!

- [0] Balbus, S. A. and Papaloizou, J. C. B. (1999).
On the Dynamical Foundations of α Disks.
Astrophysical Journal, 521:650–658.
- [0] Just, A., Yurin, D., Makukov, M., Berczik, P., Omarov, C.,
Spurzem, R., and Vilkoviskij, E. Y. (2012).
Enhanced Accretion Rates of Stars on Supermassive Black
Holes by Star-Disk Interactions in Galactic Nuclei.
The Astrophysical Journal, 758:51.
- [0] Kennedy, G. F., Meiron, Y., Shukirgaliyev, B., Panamarev, T.,
Berczik, P., Just, A., and Spurzem, R. (2016).
Star-disc interaction in galactic nuclei: orbits and rates of
accreted stars.
Monthly Notices of the Royal Astronomical Society,
460:240–255.

- [0] Kuiper, R., Klahr, H., Beuther, H., and Henning, T. (2010).
Circumventing the Radiation Pressure Barrier in the Formation
of Massive Stars via Disk Accretion.
Astrophysical Journal, 722:1556–1576.
- [0] Lodato, G. (2008).
Classical disc physics.
New Astronomy Review, 52:21–41.
- [0] Mignone, A., Bodo, G., Massaglia, S., Matsakos, T., Tesileanu,
O., Zanni, C., and Ferrari, A. (2007).
PLUTO: A Numerical Code for Computational Astrophysics.
The Astrophysical Journal Supplement Series, 170:228–242.
- [0] Novikov, I. D. and Thorne, K. S. (1973).
Astrophysics of black holes.
In Dewitt, C. and Dewitt, B. S., editors, *Black Holes (Les
Astres Occlus)*, page 343–450.

- [0] Shakura, N. I. and Sunyaev, R. A. (1973).
Black holes in binary systems. Observational appearance.
Astronomy & Astrophysics, 24:337–355.